

# Running TESLA on the Z-pole

Klaus Mönig

DESY-Zeuthen

- Introduction
- Measurements of electroweak quantities on the Z
- B-physics on the Z
- Conclusions

LEP+SLD+TEVATRON measure electroweak observables on the permille level

Quantities:

- **Z-lineshape:** Partial widths of  $Z \rightarrow f\bar{f}$ ,  $\Delta\rho$ ,  $N_\nu$
- **Asymmetries:** Weak mixing angle in Z-decays,  $\sin^2\theta_{\text{eff}}^l$
- **b-quark partial width and asymmetries** ( $R_b$ ,  $\mathcal{A}_b$ )  
Mass dependent vertex corrections
- **W-mass:**  $\Delta r$

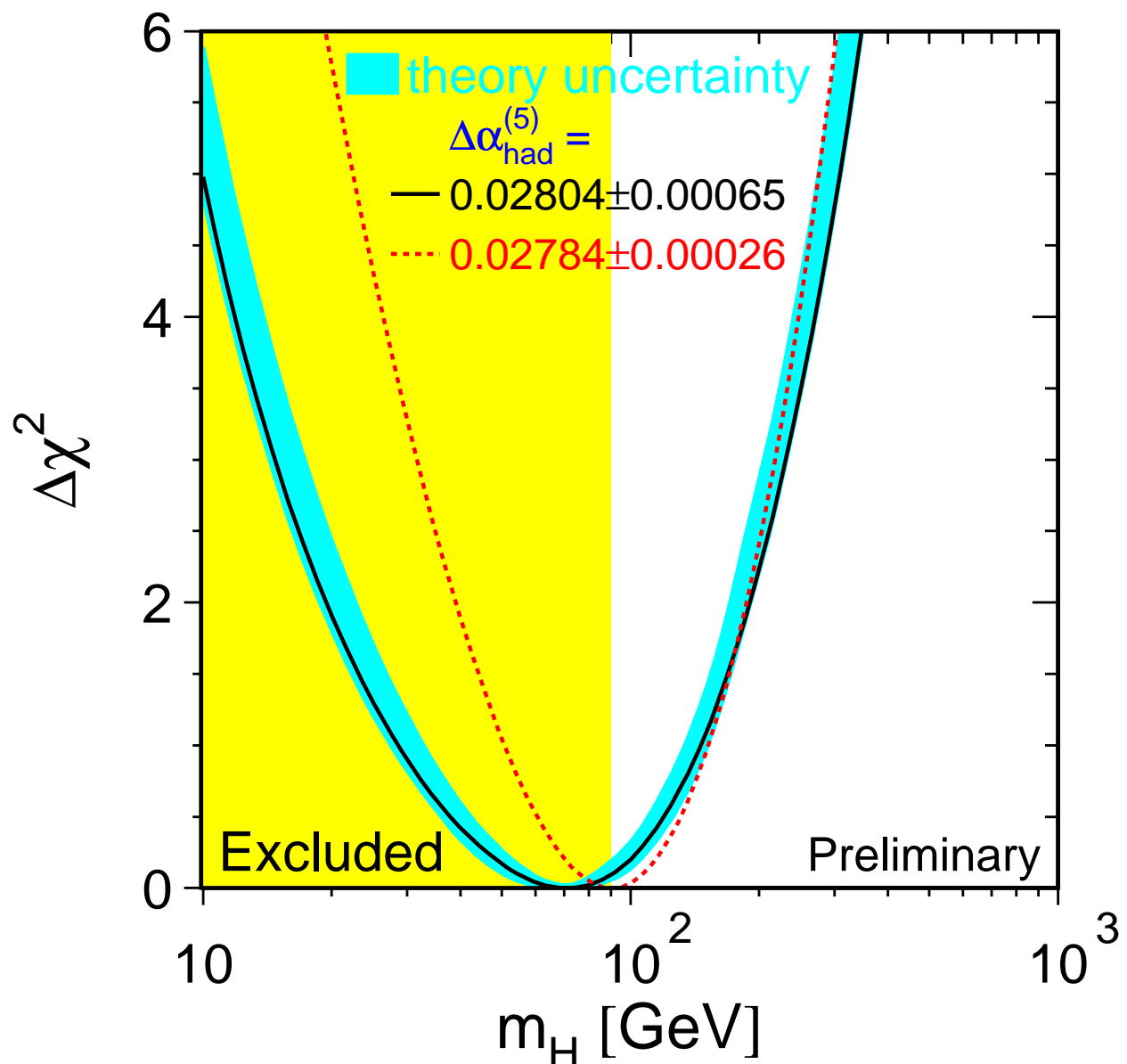
Present situation:

- **LEP:**  $\sim 4 \times 4 \cdot 10^6$  Zs with unpolarised beams  
 $\sim 4 \times 200\text{pb}^{-1}$  above the W-threshold
- **SLD:**  $\sim 5.5 \cdot 10^5$  Zs with  $\mathcal{P} \sim 75\%$  electron polarisation

## Objectives

- Test consistency of the theory on the loop level
- Try to “measure” unknown parameters in the theory

E.g. Higgs mass sensitivity of current precision data:



- The linear collider can produce  $\sim 10^9$  Zs on resonance (corresponds to  $\sim 30\text{fb}^{-1}$  or 50 days)  
 $\mathcal{L} = 7 \cdot 10^{33} \text{cm}^{-2}\text{s}^{-1} \Rightarrow 230 \text{ Hz of } Z \rightarrow q\bar{q}$
- electrons and positrons can be polarised with  $\mathcal{P}_{e^-} = \pm 80\%$ ,  $\mathcal{P}_{e^+} = \pm 60\%$   
 (corresponds to an effective polarisation of  $\frac{\mathcal{P}_{e^+} + \mathcal{P}_{e^-}}{1 + \mathcal{P}_{e^+}\mathcal{P}_{e^-}} \sim 95\%$ )
- positive and negative polarisations can be switched randomly from bunch to bunch (or train to train) independent for electrons and positrons
- polarimeters are available for relative measurements

## LEP lineshape parameters:

$$m_Z, \Gamma_Z, \sigma_0 \propto \frac{\Gamma_\ell \Gamma_{\text{had}}}{\Gamma_Z^2}, \quad R_\ell = \frac{\Gamma_{\text{had}}}{\Gamma_\ell}$$

- $m_Z$ : probably taken as reference for  $E_b$
- $\Gamma_Z$ :  $\Delta\Gamma_Z(\text{LEP}) = 2.4 \text{ MeV}$  from statistics and  $E_b$ , where  $E_b$  error is due to time dependences  
 $\Rightarrow$  With a good spectrometer to measure  $E_b$ , this can go to 1 MeV
- $\sigma_0$ : dominated by theoretical error on luminosity  
 $\Rightarrow$  assume no improvement
- $R_\ell$ : Close to systematics dominated  
 Lepton systematics is of statistical nature  
 $\rightarrow$  assume 0.05% on  $R_\ell$  from hadrons (LEP)  
 $\Rightarrow \sim 70\%$  improvement

## Physics quantities derived from the lineshape

	LEP	TESLA
$m_Z$	$91.1867 \pm 0.0021 \text{ GeV}$	$\pm 0.0021 \text{ GeV}$
$\alpha_s(m_Z)$	$0.1212 \pm 0.0034$	$\pm 0.0016$
$\Delta\rho$	$(0.42 \pm 0.12) \cdot 10^{-2}$	$\pm 0.05 \cdot 10^{-2}$
$N_\nu$	$2.994 \pm 0.0011$	$\pm 0.0011$

(For  $\alpha_s$  and  $\Delta\rho$   $N_\nu = 3$  is assumed)

- $\alpha_s$  and  $\Delta\rho$  give an interesting improvement
- However did not exploit (yet) the ability of the better detector to classify events
- One should also think if lower statistics processes like large angle Bhabha or  $e^+e^- \rightarrow \gamma\gamma$  could give a theoretically better understood luminosity determination

scale DELPHI analysis:

$$\begin{aligned} R_b &= 0.21634 \pm 0.00075 \text{ (stat dat + MC)} \\ &\pm 0.00028 \text{ (uds - bg)} \\ &\pm 0.00030 \text{ (c - bg)} \\ &\pm 0.00027 \text{ (hem corr)} \end{aligned}$$

DELPHI working point:  $\varepsilon_b \approx 30\%$  purity  $\approx 98\%$   
Possible for TESLA:  $\varepsilon_b \approx 40\%$  purity  $\approx 99.5\%$

- statistical error down by a factor 20
- c-background down by a factor 4
- uds-background mainly from gluon splitting to  $b\bar{b}$  can be measured much better with TESLA
- hemisphere correlation is mainly QCD
  - detector resolution factor 10 better than LEP
  - losses are mainly due to mass cut (Lorenz invariant)
  - energy dependence should be much smaller
  - also this source should decrease by a factor 4-5
- $\Delta R_b = 0.00014$  should be possible (factor 5 to LEP)

## Definition

$$\sigma = \sigma_u [1 - \mathcal{P}_{e^+} \mathcal{P}_{e^-} + A_{LR} (\mathcal{P}_{e^+} - \mathcal{P}_{e^-})]$$

with  $\mathcal{P}_{e^+}$  ( $\mathcal{P}_{e^-}$ ) longitudinal polarisations of the positrons (electrons)

$A_{LR}$  measures weak mixing angle  $\sin^2 \theta_{\text{eff}}^l$ :

$$A_{LR} = \mathcal{A}_l$$

$$\mathcal{A}_l = \frac{2g_{Vl}g_{Al}}{g_{Vl}^2 + g_{Al}^2}$$

$$\frac{g_{Vl}}{g_{Al}} = 1 - 4|Q_l| \sin^2 \theta_{\text{eff}}^l$$

- $\sin^2 \theta_{\text{eff}}^l$  is a very sensitive variable to see loop corrections to the Z-couplings.
- $A_{LR}$  is the variable most sensitive to  $\sin^2 \theta_{\text{eff}}^l$



Four independent measurements:

(4 combinations with positive/negative electron/positron polarisation)

$$\sigma_{++} = \sigma_u [1 - \mathcal{P}_{e^+}\mathcal{P}_{e^-} + A_{\text{LR}}(\mathcal{P}_{e^+} - \mathcal{P}_{e^-})]$$

$$\sigma_{-+} = \sigma_u [1 + \mathcal{P}_{e^+}\mathcal{P}_{e^-} + A_{\text{LR}}(-\mathcal{P}_{e^+} - \mathcal{P}_{e^-})]$$

$$\sigma_{+-} = \sigma_u [1 + \mathcal{P}_{e^+}\mathcal{P}_{e^-} + A_{\text{LR}}(\mathcal{P}_{e^+} + \mathcal{P}_{e^-})]$$

$$\sigma_{--} = \sigma_u [1 - \mathcal{P}_{e^+}\mathcal{P}_{e^-} + A_{\text{LR}}(-\mathcal{P}_{e^+} + \mathcal{P}_{e^-})]$$

$\implies A_{\text{LR}}$  can be measured without knowing  $\mathcal{P}_{e^+}, \mathcal{P}_{e^-}$ :

$$A_{\text{LR}} = \sqrt{\frac{(\sigma_{++} + \sigma_{--} - \sigma_{+-} - \sigma_{-+})(-\sigma_{++} + \sigma_{--} - \sigma_{+-} + \sigma_{-+})}{(\sigma_{++} + \sigma_{--} + \sigma_{+-} + \sigma_{-+})(-\sigma_{++} + \sigma_{--} + \sigma_{+-} - \sigma_{-+})}}$$

About 10% of the statistics is needed on the small cross sections

Only difference between  $|\mathcal{P}_{e^\pm}^+|$  and  $|\mathcal{P}_{e^\pm}^-|$  needs to be known from polarimetry

Can be brought under control with polarimeters a la SLD

Polarisation difference ( $\Delta\mathcal{P}_{e^\pm} = |\mathcal{P}_{e^\pm}^+| - |\mathcal{P}_{e^\pm}^-|$ ):

- Need SLD like polarimeter
- Asymmetry in one polarimeter channel:  
 $A_i = a_i \mathcal{P}_e \mathcal{P}_\gamma$  ( $a_i$  = analysing power)
- Laser polarisation can be switched pulse to pulse
- Allow for different laser currents dependent on the polarisation
- Need two polarimeter channels with different analysing power
- combined fit of Z-rates and polarimeter rates can get  $\Delta\mathcal{P}_{e^\pm}$  and  $a_i$  as well
- However need polarimeter counting rates about 10 times the Z rate (ok for SLD)

## Statistical precision:

$$\Delta A_{\text{LR}} = 4 \cdot 10^{-5} \cdot \sqrt{\frac{10^9}{N_Z}}$$

## Systematic uncertainties

- **Beam energy:**  $\Delta A_{\text{LR}}/\Delta\sqrt{s} \approx 2 \cdot 10^{-2} / \text{GeV}$   
 $\Rightarrow$  need  $\Delta\sqrt{s} \approx 1 \text{ MeV}$
- **Luminosity difference:** Only relative precision needed.  
Should be no problem if luminometer inside the mask is possible
- **Backgrounds:** To be kept below  $10^{-4}$   
According to LEP experience no problem
- **Beamstrahlung:**  $\Delta A_{\text{LR}} = 9 \cdot 10^{-4}$   
Needs to be known on the few percent level

Assume  $\Delta A_{\text{LR}} = 10^{-4} \Rightarrow \Delta \sin^2 \theta_{\text{eff}}^{\ell} = 0.000013$

Without polarised beams (LEP) the forward-backward asymmetries can be measured:

$$\begin{aligned} A_{FB}^q &= \frac{\sigma_F^{(q)} - \sigma_B^{(q)}}{\sigma_T^{(q)}} \\ &= \frac{3}{4} \mathcal{A}_e \mathcal{A}_q \end{aligned}$$

With polarised beams (SLD, TESLA) the left-right-forward-backward asymmetries can be measured:

$$\begin{aligned} A_{FB,LR}^q &= \frac{\sigma_{L,F}^{(q)} - \sigma_{L,B}^{(q)} - \sigma_{R,F}^{(q)} + \sigma_{R,B}^{(q)}}{\sigma_L^{(q)} + \sigma_R^{(q)}} \\ &= \frac{3}{4} \mathcal{P} \mathcal{A}_q \end{aligned}$$

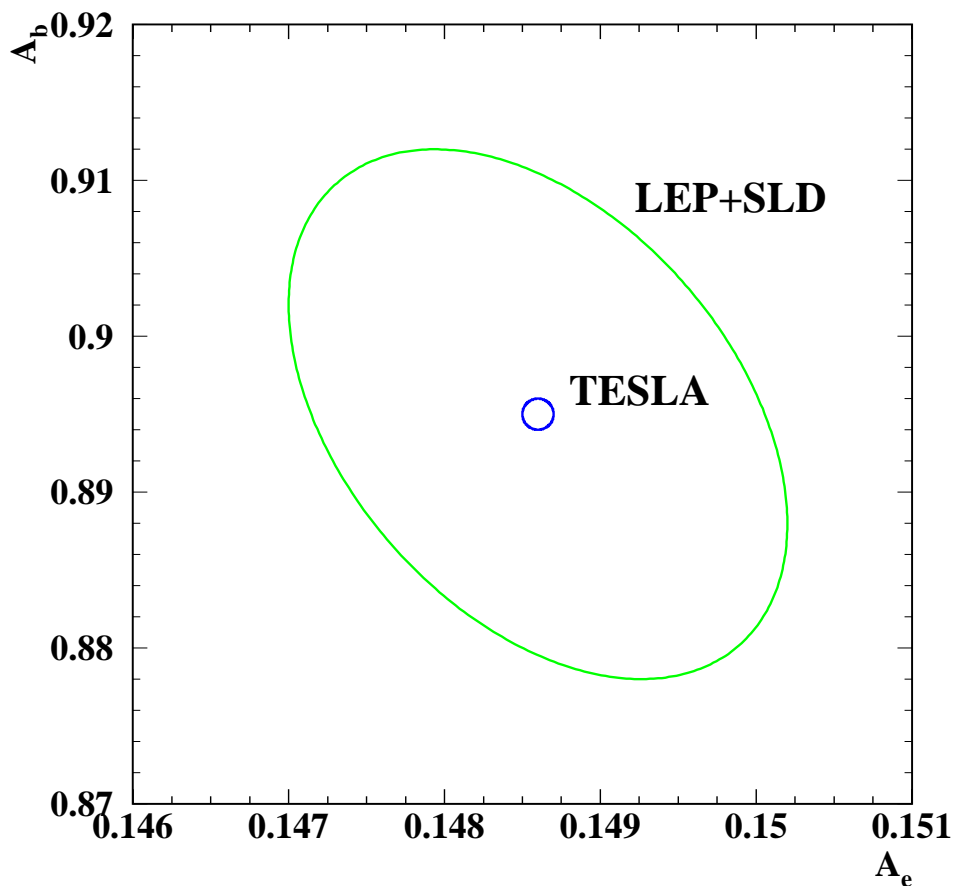
Statistically factor  $\mathcal{P}/\mathcal{A}_e \sim 6$  more sensitive to  $\mathcal{A}_b$

However most systematics scale with the asymmetry

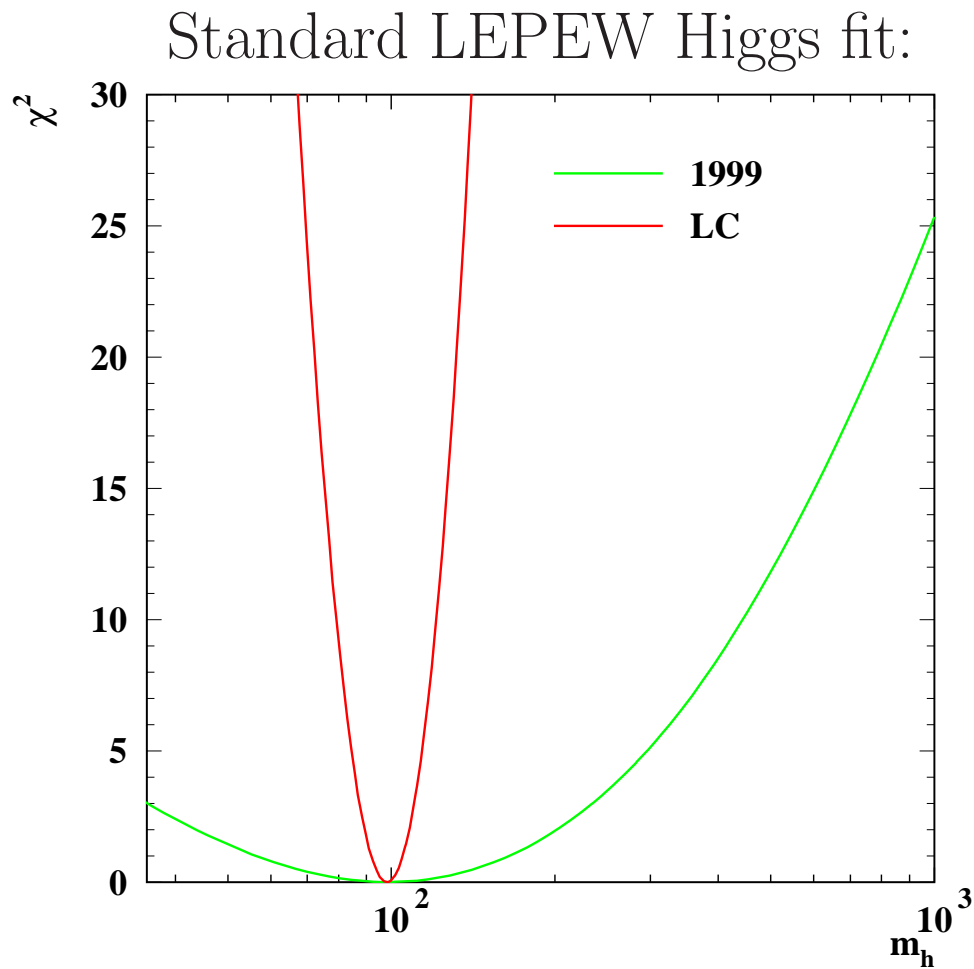
Two main techniques: leptons and jetcharge

- Statistical error  $\Delta\mathcal{A}_b \simeq 4 \cdot 10^{-4}$  in both cases
- Light quark systematics can be reduced by a (harder) lifetime tag
- For jetcharge reduce hemisphere correlations by a thrust cut
- leptons will be dominated by  $B\bar{B}$ -mixing (statistical error!)
- A total error of  $\Delta\mathcal{A}_b = 1 \cdot 10^{-3}$  seems realistic

Similar improvement as for  $\mathcal{A}_e$

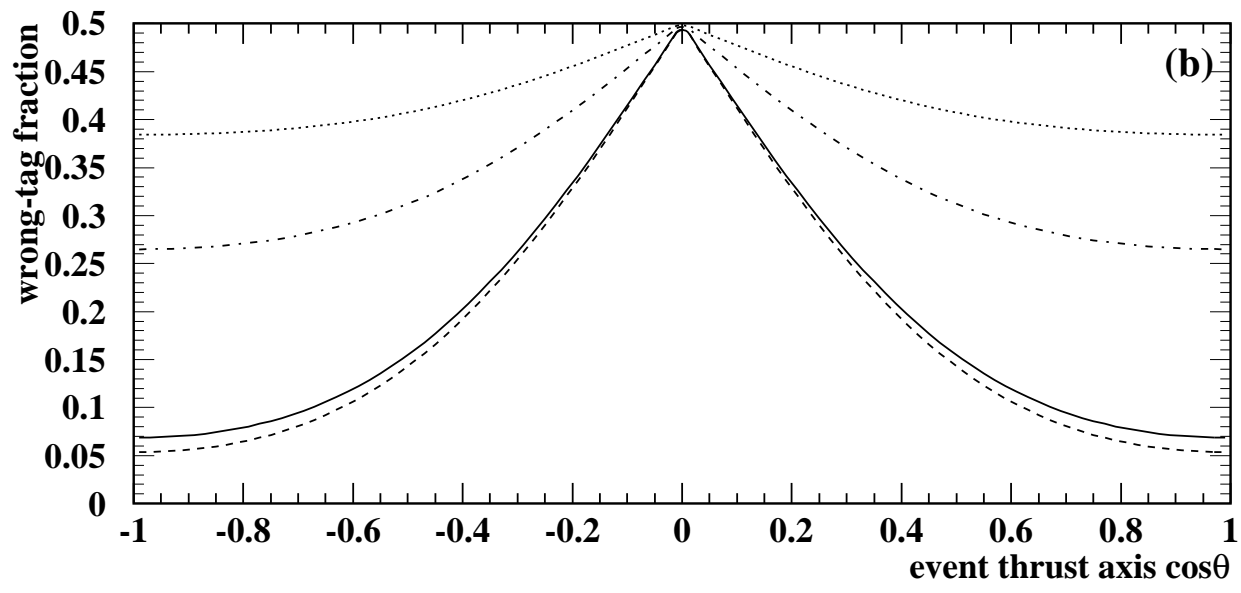
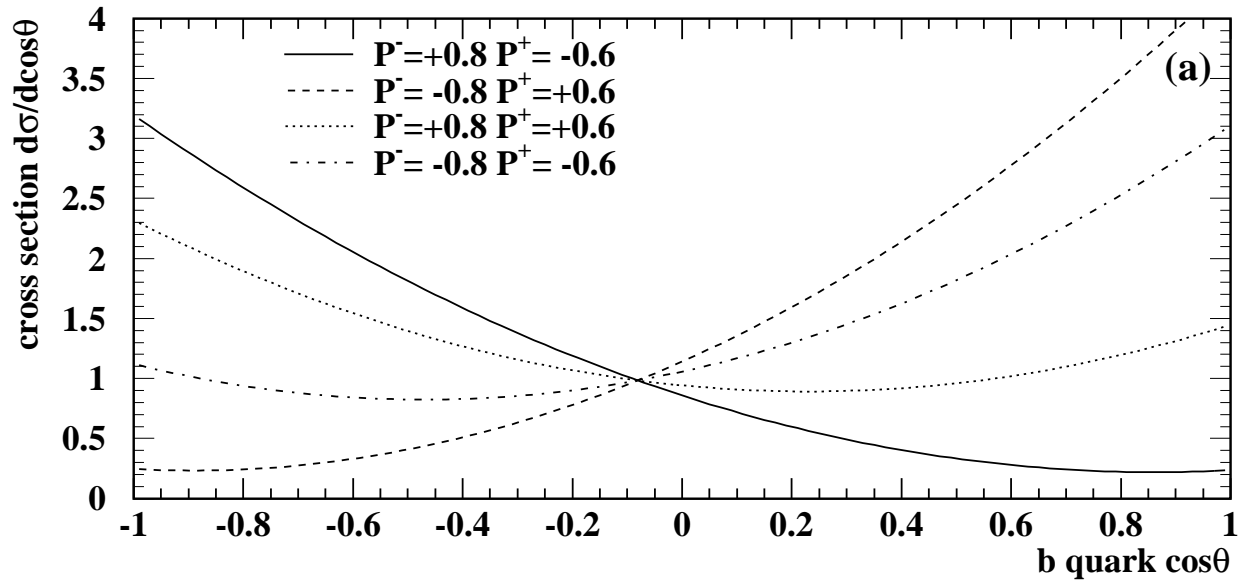


If no new physics found up to then:

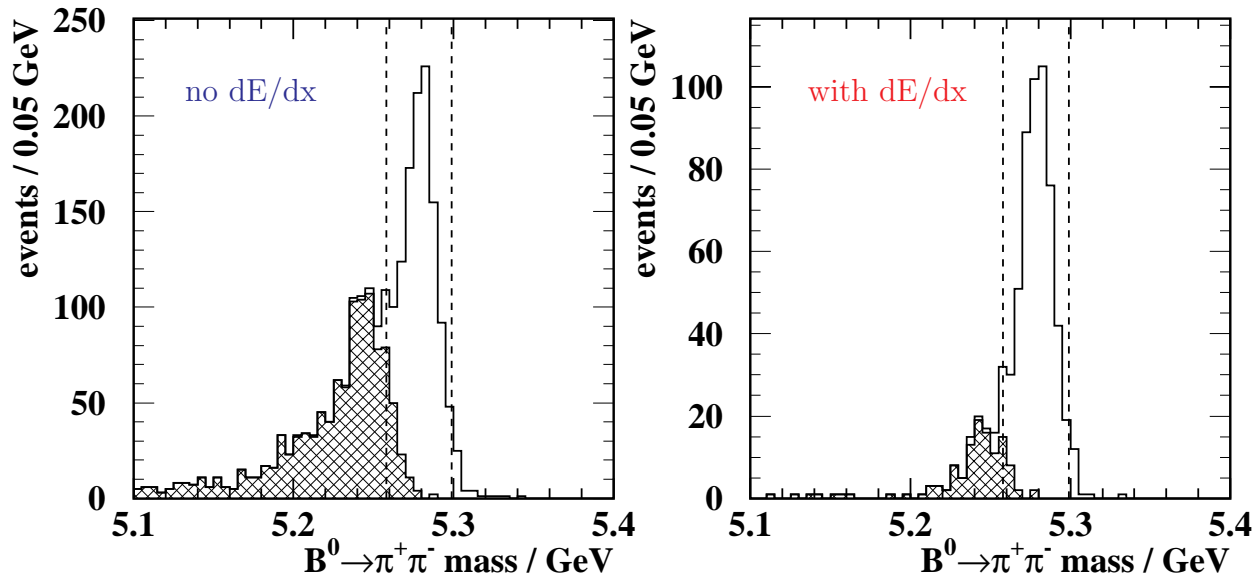


More theory:  $\rightarrow$  S. Heinemeyers talk

- $4 \cdot 10^8$  b-hadrons
- Polarisation give primary flavour tagging “for free”



- Missing particle ID can be replaced by excellent momentum resolution



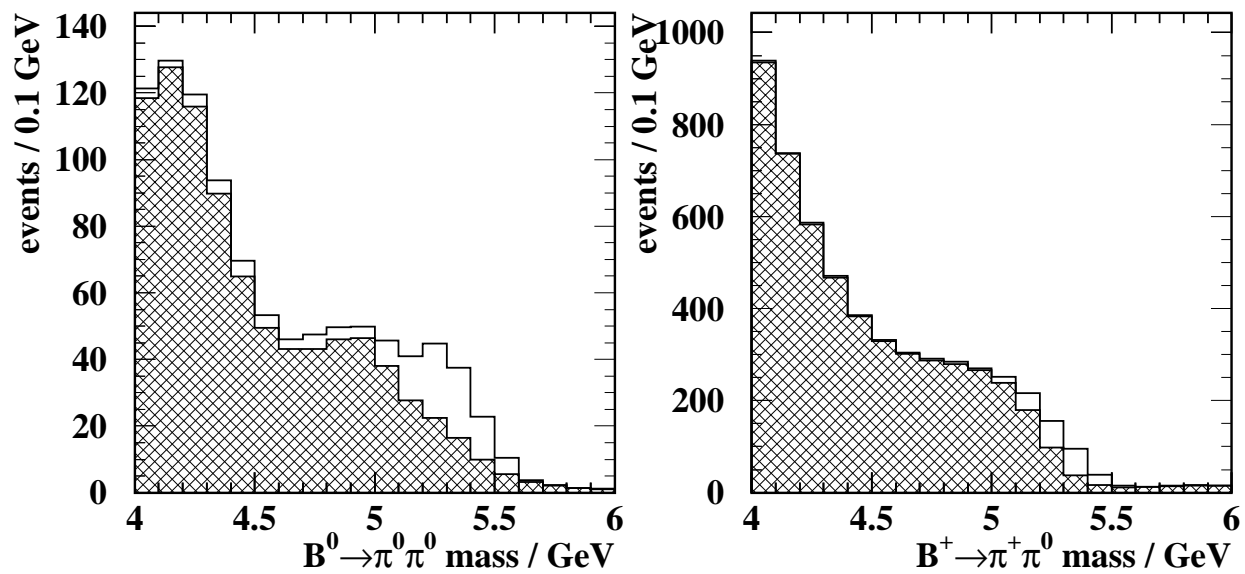
## Results

	$\sin 2\beta$	“ $\sin 2\alpha$ ”
BarBar	0.12	0.26
CDF	0.08	0.10
ATLAS	0.02	0.14
LHC-b	0.01	0.05
<b>TESLA</b>	<b>0.04</b>	<b>0.07</b>

Not the best, but interesting cross check!



- needed to disentangle direct from penguin contributions in  $B^0 \rightarrow \pi^+ \pi^-$
- only possible in  $e^+ e^-$ -machines
- can be done at LC with good calorimetry and good b-tagging/anti-b-tagging



- Competitive results to BarBar can be obtained

- With less than a year of running on the Z huge progress on  $\sin^2 \theta_{eff}^l$  can be made
- Also Zbb-couplings can be improved a lot
- Only modest improvement in Z-lineshape observables possible
- Some interesting cross checks in B-physics, however no “golden channel” (yet)