

ASSOCIATED PRODUCTION

OF Higgs WITH SQUARKS (\tilde{t})

at future e^+e^- colliders

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with $\left\{ \begin{array}{l} \text{A. DJOUADI} \\ \text{G. MOULTAKA} \end{array} \right.$

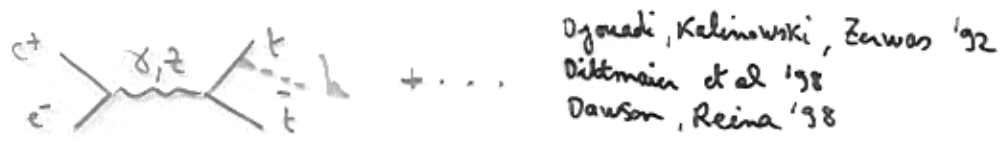
hep-ph/9903218

"Physics and Experiments with future linear e^+e^- colliders"

Sitges April '95

MOTIVATIONS / SUMMARY :

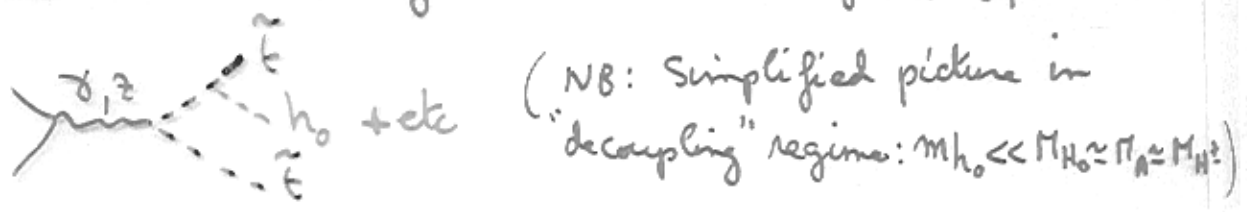
SM Higgs production in association with $t\bar{t}$ [in pp or e^+e^-]:



NOT dominant % $e^+e^- \rightarrow z \rightarrow zh, ww \rightarrow h$;

BUT further background reduction possibilities
(lepton tagging $t \rightarrow wb \rightarrow l\nu$)

•• MSSM "analogous": associated (lighter) h_0 production:



MAYBE A PRIORI COMPETITIVE (DOMINANT?)
especially for large (maximal?) $\tilde{\theta}_t$ mixing scenario

$$\begin{pmatrix} M_{QL}^2 + D + m_t^2 & m_t(A_t - \mu/\tan\beta) \\ m_t(A_t - \mu/\tan\beta) & M_{tR}^2 + D + m_t^2 \end{pmatrix} \rightarrow m_{\tilde{t}_1}^2 \ll m_{\tilde{t}_2}^2 \lesssim m_t \quad \left| \begin{array}{l} \text{for } \tilde{\theta}_t \text{ "max"} \\ |\tilde{\theta}_t| \approx \frac{\pi}{4} \end{array} \right.$$

→ obvious phase space enhancement of $\sigma(\tilde{t}\tilde{t}^*h)$;
but also because $g_{\tilde{t}_1\tilde{t}_1 h_0} \uparrow$ for $(A_t - \mu/\tan\beta)$ large enough
 $\equiv \tilde{A}_t$

→ pp study quite promising (Djouadi, Maultaka, JLK
PRL 80 (1998)
hep-ph/9711244)
WHAT ABOUT e^+e^- ?

$\phi\phi, \phi\phi\phi, \phi\phi\phi\phi$ F, D terms + SOFT-BREAKING TERMS

3)

$$\rightarrow M_{\tilde{t}}^2 = \begin{pmatrix} M_{\tilde{Q}_L}^2 + m_t^2 + (\frac{1}{2} - \frac{2}{3} s_w^2) m_t^2 \cos 2\beta & m_t (A_t - \mu/\tan\beta) \\ m_t (A_t - \mu/\tan\beta) & M_{\tilde{t}_R}^2 + m_t^2 + \frac{2}{3} s_w^2 m_t^2 \cos 2\beta \end{pmatrix}$$

$$\rightarrow m_{\tilde{t}_{1,2}}^2 = m_t^2 + \frac{1}{2} \left[M_{\tilde{Q}_L}^2 + M_{\tilde{t}_R}^2 + \frac{m_t^2}{2} \cos 2\beta \pm \sqrt{\left(M_{\tilde{Q}_L}^2 - M_{\tilde{t}_R}^2 + (\frac{1}{2} - \frac{2}{3} s_w^2) m_t^2 \cos 2\beta \right)^2 + 4 m_t^2 (A_t - \mu/\tan\beta)^2} \right]$$

$$\tilde{\theta}_t \equiv \frac{1}{2} \text{Atan} \left[\frac{2 m_t (A_t - \mu/\tan\beta)}{(M_{\tilde{Q}_L}^2 - M_{\tilde{t}_R}^2) + D..} \right] \quad (\tan\beta = \frac{v_u}{v_d})$$

and physical $h_{\tilde{t}_1, \tilde{t}_2}$, etc couplings obtained from

$$\begin{pmatrix} \tilde{t}_L \\ \tilde{t}_R \end{pmatrix} = \begin{pmatrix} \cos \tilde{\theta}_t & -\sin \tilde{\theta}_t \\ \sin \tilde{\theta}_t & \cos \tilde{\theta}_t \end{pmatrix} \begin{pmatrix} \tilde{t}_1 \\ \tilde{t}_2 \end{pmatrix} ; \begin{pmatrix} H_d \\ H_u \end{pmatrix} = \begin{pmatrix} \cos \alpha & -\sin \alpha \\ \sin \alpha & \cos \alpha \end{pmatrix} \begin{pmatrix} H \\ h \end{pmatrix}$$

\rightarrow

$$g_{\tilde{t}_1, \tilde{t}_2, h}^{\tilde{t}_1, \tilde{t}_2} = -\sin(\alpha + \beta) \left[\cos^2 \tilde{\theta}_t - \frac{2}{3} s_w^2 \cos(2\tilde{\theta}_t) \right] + \frac{m_t^2 \cos \alpha}{m_z^2 \sin \beta} + \frac{\sin(2\tilde{\theta}_t)}{2} \frac{m_t}{m_z^2} \left(A_t \frac{C_\alpha}{S_\beta} + \mu \frac{S_\alpha}{S_\beta} \right)$$

$$g_{\tilde{t}_2, \tilde{t}_2, h}^{\tilde{t}_2, \tilde{t}_2} = -\sin(\alpha + \beta) \left[\sin(2\tilde{\theta}_t) \left(\frac{2}{3} s_w^2 - \frac{1}{4} \right) \right] + \frac{\cos(2\tilde{\theta}_t)}{2} \frac{m_t}{m_z^2} \left(A_t \frac{C_\alpha}{S_\beta} + \mu \frac{S_\alpha}{S_\beta} \right)$$

$$\times \left(\frac{M_w}{C_w} a g \right)$$

- Decoupling regime: $m_h \ll M_H = M_A \approx M_{H^\pm} \rightarrow \alpha \rightarrow \beta - \frac{\pi}{2}$

$$\text{So } g_{\tilde{t}_1, \tilde{t}_1, h}^{\tilde{t}_1, \tilde{t}_1} \sim \frac{\cos 2\beta}{2} \left[\cos^2 \tilde{\theta}_t - \frac{2}{3} s_w^2 \cos(2\tilde{\theta}_t) \right] + \frac{m_t^2}{m_z^2} + \frac{m_t \sin(2\tilde{\theta}_t)}{m_z^2} (A_t - \mu/\tan\beta)$$

originally $\tilde{t}_L \tilde{t}_R H = \cos \tilde{\theta}_t \tilde{t}_1 \tilde{t}_2 H + \dots$

Now: "large" $\tilde{A}_t (\equiv A_t - \mu/\tan\beta)$ limit?

clearly, $m_{\tilde{t}_1} \rightarrow$ while $g_{\tilde{t}_1, \tilde{t}_1, h}^{\tilde{t}_1, \tilde{t}_1} \sim \sin(2\tilde{\theta}_t) \tilde{A}_t \sim -\tilde{A}_t$

($g_{\tilde{t}_1, \tilde{t}_2}^{\tilde{t}_1, \tilde{t}_2} \sim \cos(2\tilde{\theta}_t) \tilde{A}_t$) $\xrightarrow{\beta \rightarrow \frac{\pi}{4}}$ linear growth!

MAINLY: $e^+e^- (\gamma\gamma) \rightarrow \tilde{t}_1 \tilde{t}_1^* h$

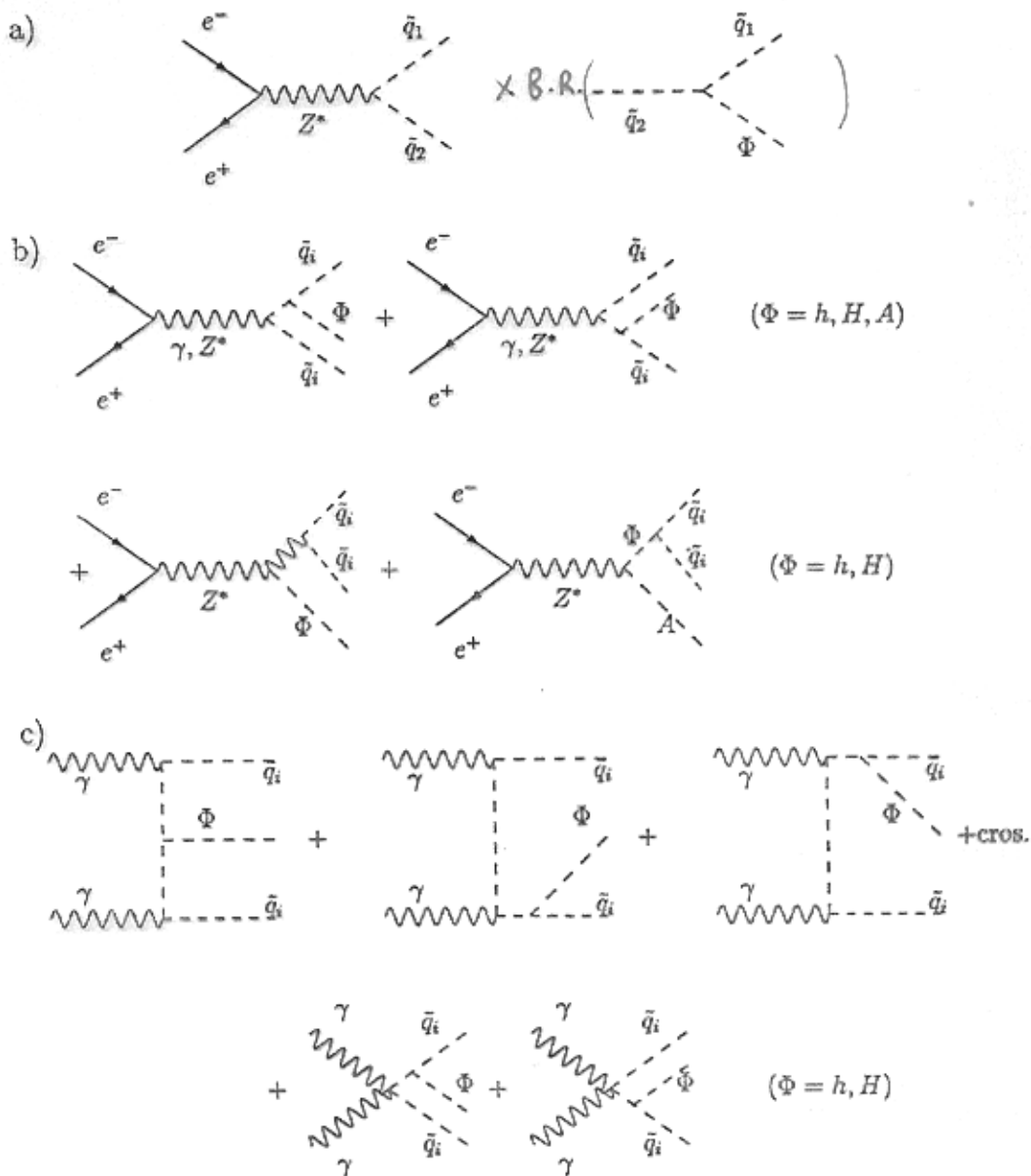


Figure 1: Feynman diagrams for the production of a Higgs boson Φ in association with a pair of squarks in e^+e^- and $\gamma\gamma$ collisions.

σ [86]

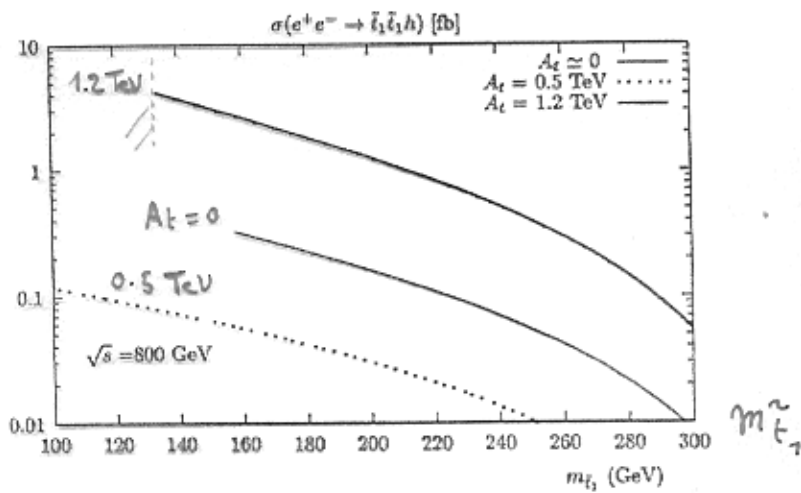


Figure 3: The cross section $\sigma(e^+e^- \rightarrow \tilde{t}_1 \tilde{t}_1 h)$ [in fb] as a function of the \tilde{t}_1 mass and three different choices of the other parameters: $\tan \beta = 30$, $A_t = 0$ (0.5) TeV; $\tan \beta = 3$, $A_t = 1.2$ TeV.

HOWEVER, TH. + EXP. CONSTRAINTS on large A_t :

1) Charge & Color non-breaking constraint(s) :

$$A_t^2 < 3(m_{QL}^2 + m_{tR}^2 + \mu^2 + m_{\tilde{E}U}^2) \quad (\text{at EWSB scale})$$

2) The (in)famous most "new physics killer" :

$$\Delta\rho = \epsilon_1 \approx \text{W loop} - \text{Z loop} \quad [\text{mainly}; + \text{also other oblique corrections}]$$

measuring $SU(2)$ "custodial" breaking

MS: large $(m_t^2 - m_b^2) \times G_F$ contribution

MSSM: additional $SU(2)$ -breaking sources!

$$\bullet \tilde{m}_{tL} \tilde{m}_{tR} \neq \tilde{m}_{bR}$$

$$\bullet m_t \tilde{A}_t \neq m_b \tilde{A}_b$$

$$\Delta\rho_{\text{MSSM}}^{(\tilde{t}, \tilde{b})} \approx \frac{3 G_F}{8\pi^2 \sqrt{2}} \left[c_{\tilde{\theta}_t}^2 (c_{\tilde{\theta}_b}^2 |m_{\tilde{t}_1}^2 - m_{\tilde{b}_1}^2| + s_{\tilde{\theta}_b}^2 |m_{\tilde{t}_1}^2 - m_{\tilde{b}_2}^2|) \right. \\ \left. + s_{\tilde{\theta}_t}^2 (c_{\tilde{\theta}_b}^2 |m_{\tilde{t}_2}^2 - m_{\tilde{b}_1}^2| + s_{\tilde{\theta}_b}^2 |m_{\tilde{t}_2}^2 - m_{\tilde{b}_2}^2|) \right. \\ \left. - c_{\tilde{\theta}_t}^2 s_{\tilde{\theta}_t}^2 |m_{\tilde{t}_1}^2 - m_{\tilde{t}_2}^2| - c_{\tilde{\theta}_b}^2 s_{\tilde{\theta}_b}^2 |m_{\tilde{b}_1}^2 - m_{\tilde{b}_2}^2| \right] \\ \propto \sin^2(2\theta)$$

Barbieri, Petriani '83

Drees, Hagihara '90

$$\delta(\Delta\rho) \gtrsim 2-3 \cdot 10^{-3} (1.2\sigma) \quad \text{e.g. Altarelli's reviews '97 '99}$$

3) + direct limits:

$$m_h \gtrsim 50 \text{ GeV} \quad (\text{decoupling regime}) \quad (\text{LEP2})$$

$$m_{\tilde{t}_1} \gtrsim 80 - 120 \text{ GeV} \quad (\text{LEP2, CDF})$$

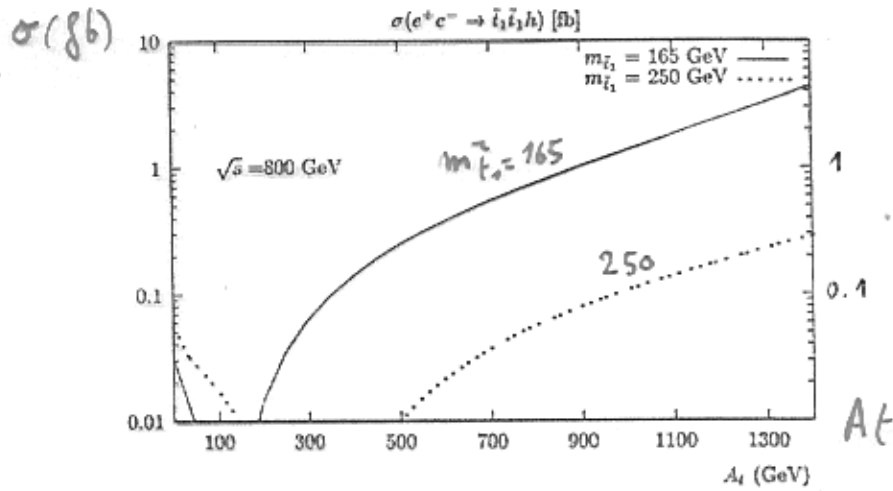


Figure 4: The cross section $\sigma(e^+e^- \rightarrow \tilde{t}_1 \bar{\tilde{t}}_1 h)$ [in fb] as a function of A_t for fixed $m_{\tilde{t}_1} = 165$ ($\tan \beta = 3$); and 250 GeV ($\tan \beta = 30$).

NB : to compare, $\sigma(t\bar{t}h)_{SM} \approx 2 \text{ fb}$ ($M_h \approx 130 \text{ GeV}$)

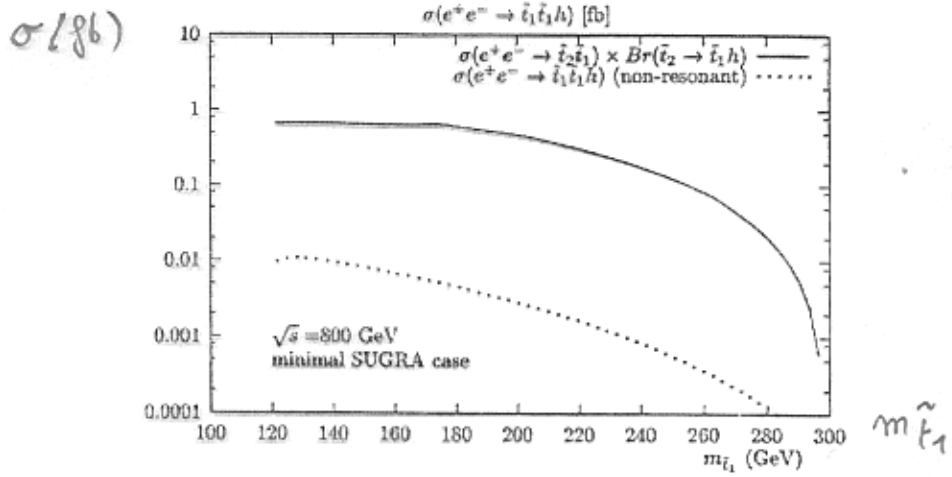


Figure 2: The production cross section $\sigma(e^+e^- \rightarrow \tilde{t}_1 \tilde{t}_1 h)$ [in fb] as a function of $m_{\tilde{t}_1}$ in the mSUGRA case; $\tan\beta = 30$, $m_{1/2} = 100$ GeV and $A_0 = -600$ GeV.

(i.e. moderate mixing case)

NB: $\tilde{t}_1 \tilde{t}_2 \sim \sin(2\tilde{\theta}_t)$

$$\tilde{t}_2 \tilde{t}_1 h \approx \cos 2\beta \left(\frac{2}{3} \sin^2 \frac{1}{4} \right) \sin(2\theta) + \frac{m_t}{2m_{\tilde{t}_2}} \tilde{A}_t \cos(2\tilde{\theta}_t)$$

can't be too large in most of the parameter space

$\Rightarrow \sigma^{\text{resonant}} \approx \sin 2\tilde{\theta}_t \cos 2\tilde{\theta}_t$

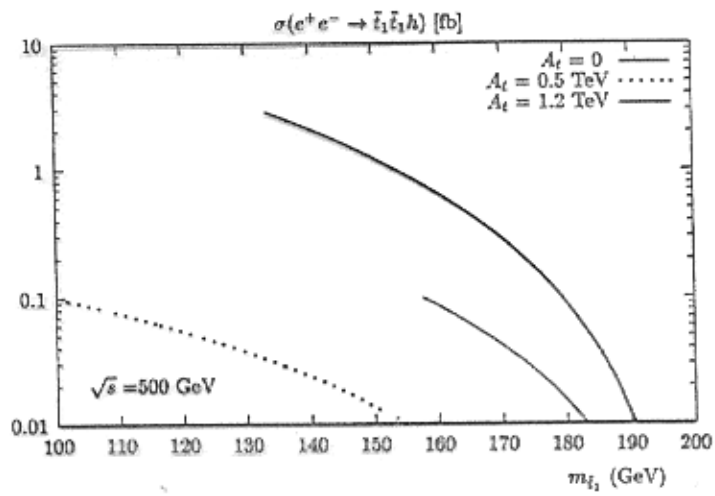


Figure 5: Same as Fig. 8 but with $\sqrt{s} = 500 \text{ GeV}$.

$\gamma\gamma$ mode

$\sigma(\text{fb})$

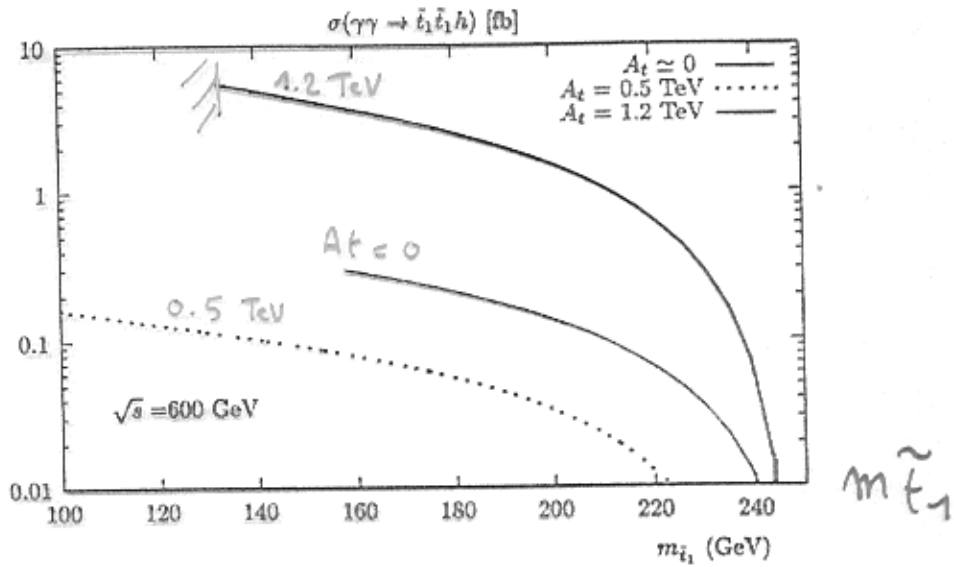


Figure 6: The cross section $\sigma(\gamma\gamma \rightarrow \tilde{t}_1 \tilde{t}_1 h)$ in $\gamma\gamma$ collisions [in fb] at a center of mass energy $\sqrt{s_{\gamma\gamma}} = 600$ GeV as a function of $m_{\tilde{t}_1}$. The other parameters have the same values as in Fig. 10.

Decay properties - Possible detection modes

1. $\tilde{t}_1 \rightarrow b\chi^+$ (if $\tilde{m}_{t_1} < m_t + m_{\chi_0}$)

2. $\tilde{t}_1 \rightarrow t\chi_0^1$ (LSP) otherwise

(also $\tilde{t}_1 \rightarrow c\chi_0$ may compete if $m_{\tilde{t}_1} - m_{\chi_0^+} \approx 0$.
(1-loop))

$\Rightarrow \tilde{t}_1 \rightarrow b\chi^+$ dominates where $\sigma_{\tilde{t}_1\tilde{t}_1 h}$ is large
 $\hookrightarrow \chi^0 W^+$
 $\hookrightarrow e\nu$

$$\tilde{t}_1 \rightarrow bW^+ + \cancel{E}$$

same topology as $t \rightarrow bW^+$

BUT \neq large amount of \cancel{E}

$ch \rightarrow b\bar{b}$ dominant decay

$\rightarrow 4b + 2W$ final state + \cancel{E}

if efficient microvertex detectors (b tagging)

should be relatively easy in clean e^+e^- environment

CONCLUSIONS

- new h production process
- $\sigma(e^+e^-, \gamma\gamma \rightarrow \tilde{q}; \tilde{q}; \phi)$ general expressions available
(rather simple analytical formulae)
- at e^+e^- LC, $\sqrt{s} = 500 - 800$ GeV
with $\int \mathcal{L} dt \sim 100 - 500 \text{ fb}^{-1}$
several hundred of events possible if $m_{\tilde{t}} < 250$ GeV
and A_t large enough
- Rather spectacular and clean signal
 $4b's + 2W's + \cancel{E}$
- Also more th. motivations:
a possible handle to probe
the mysterious SOFT-SUSY-BREAKING sector
(if A_t large, $\tilde{F}\tilde{E}h$ potentially largest MSSM coupling!)